

# FTIR Spectroscopy of Secondary-Structure Reorientation of Melibiose Permease Modulated by Substrate Binding

Natàlia Dave,\* Víctor A. Lórenz-Fonfría,\* Gérard Leblanc,<sup>†</sup> and Esteve Padrós\*

\*Unitat de Biofísica, Departament de Bioquímica i de Biologia Molecular, Facultat de Medicina, and Centre d'Estudis en Biofísica, Universitat Autònoma de Barcelona, Barcelona, Spain; and <sup>†</sup>Institut de Biologie et Technologies-Saclay, Service de Bioénergétique, Biologie Structurale et Mécanismes, Gif sur Yvette, France

**ABSTRACT** Analysis of infrared polarized absorbance spectra and linear dichroism spectra of reconstituted melibiose permease from *Escherichia coli* shows that the oriented structures correspond mainly to tilted transmembrane  $\alpha$ -helices, forming an average angle of  $\sim 26^\circ$  with the membrane normal in substrate-free medium. Examination of the deconvoluted linear dichroism spectra in H<sub>2</sub>O and D<sub>2</sub>O makes apparent two populations of  $\alpha$ -helices differing by their tilt angle (helix types I and II). Moreover, the average helical tilt angle significantly varies upon substrate binding: it is increased upon Na<sup>+</sup> binding, whereas it decreases upon subsequent melibiose binding in the presence of Na<sup>+</sup>. In contrast, melibiose binding in the presence of H<sup>+</sup> causes virtually no change in the average tilt angle. The data also suggest that the two helix populations change their tilting and H/D exchange level in different ways depending on the bound substrate(s). Notably, cation binding essentially influences type I helices, whereas melibiose binding modifies the tilting of both helix populations.

## INTRODUCTION

Melibiose permease (MelB) of *Escherichia coli* couples the uphill transport of  $\alpha$ - or  $\beta$ -galactosides to the downhill inward movement of Na<sup>+</sup>, Li<sup>+</sup>, or H<sup>+</sup>. It is a representative member of the galactoside-pentose-hexuronide family. This integral membrane protein is encoded by *melB* gene (1) and consists of 473 amino acids, of which 70% are apolar. Although most bacterial transporters cotransport H<sup>+</sup> with the solute, MelB constitutes an exception, since it has the ability to transport sugars using Na<sup>+</sup>, Li<sup>+</sup>, or H<sup>+</sup>. Each cation binds to the same site, enhancing the affinity of the cotransported sugar, with Na<sup>+</sup> and Li<sup>+</sup> being the best activators (2,3). The currently accepted topological model, consisting of 12 transmembrane (TM) domains and N- and C-termini located in the cytoplasm, is strongly supported by immunological studies (4,5), extensive *melB-phoA* fusion analyses (5,6), and proteolytic mapping (7). Moreover, Fourier transform infrared (FTIR) studies showed a high  $\alpha$ -helix content (50%), giving support to the 12-TM segment topological model (8). Additionally, these studies and subsequent attenuated total reflection-Fourier transform infrared (ATR-FTIR) analyses suggest the existence of two different types of transmembrane  $\alpha$ -helices (8,9). Photolabeling (10) and mutagenesis studies combined with fluorescence analysis (11) have highlighted the importance of the helix IV loop 4-5 domain for MelB transport function. Furthermore, fluorescence and FTIR spectroscopies have revealed substrate-induced conformational changes at different

stages of the transport cycle (8,12–15). An electrophysiological approach revealed fast transient currents (20-ms range) triggered by either Na<sup>+</sup> or melibiose binding, resulting from movements of charged amino acids and/or a reorientation of helix dipoles (16).

One of the most important aims of protein research is to study the relation between structure and function. Though MelB 3D structure determination would be an essential step in this regard, FTIR spectroscopy in some of its different applications (secondary-structure quantification (8), H/D exchange (15), and substrate-induced difference spectroscopy (9,17)) has provided relevant information about structural changes related to MelB function. Determination of the orientation of the secondary-structure components of proteins with polarized incident light is another of the important applications of FTIR spectroscopy (18–20). The method is based on the fact that infrared light absorption is maximal when the transition dipole moment is parallel to the electric field component of the incident light beam. The dipole orientation can be assessed by measuring the spectrum intensity on changing the incident light polarization (parallel,  $\parallel$ , versus perpendicular,  $\perp$ ).

An electron crystallographic study of 2D MelB crystals (21) strongly suggest that at least 7 out of the 12 transmembrane (TM) domains (possibly  $\alpha$ -helices), display significant albeit variable tilt with respect to the membrane normal. A recent reconstitution of 3D MelB structure, although of limited resolution, confirms the presence of such tilted transmembrane helical domains (22). Moreover, tilted membrane helices are observed in all the 3D structures of secondary transporters successfully crystallized (23–29). It is important to note that changes in the tilting of transmembrane helices have been proposed, almost systematically, to be involved in the alternating-access mechanisms (30,31).

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Address reprint requests to Natàlia Dave, Unitat de Biologia Cel·lular i Molecular, Institut Municipal d'Investigació Mèdica, Universitat Pompeu Fabra, Barcelona, Spain. E-mail: ndave@imim.es.

Víctor A. Lórenz-Fonfría's present address is Department of Materials Science and Engineering, Nagoya Institute of Technology, Nagoya, Japan.

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In this study, we use polarized ATR-FTIR spectroscopy to quantitatively evaluate the average helix tilting of MelB with respect to the lipid membrane normal in proteoliposomes containing purified MelB transporters. Moreover, the possibility of independently manipulating the sodium and sugar substrates offers, additionally, the opportunity to investigate whether their interaction with MelB modifies the tilting of the helical domains of the transporter, in particular that of two hypothetical populations of  $\alpha$ -helices that were previously identified on the basis of their different FTIR signatures (8).

## MATERIALS AND METHODS

### Materials

p-Nitrophenyl  $\alpha$ -D-(6- $^3$ H) galactopyranoside ( $^3$ H]-NPG) was synthesized in our department (Institut de Biologie et Technologies-Saclay, CEA, France), under the direction of Dr. B. Rousseau. Synthesis of LAPAO ((3-laurylamido)-N,N'-(dimethylamino)propylamine oxide) was performed as previously described (21). Dodecyl maltoside (DM) was obtained from Boehringer Mannheim (Cambridge, MA), and Ni-NTA resin was from Qiagen (Hilden, Germany). SM-2 Bio-Beads were obtained from Bio-Rad (Hercules, CA). Total *E. coli* lipids (acetone/ether precipitated) were purchased from Avanti Polar Lipids (Alabaster, AL). High-purity-grade salts or chemicals (Suprapur, Merck, Whitehouse Station, NJ) were used to prepare nominally  $\text{Na}^+$ -free media containing  $<20 \mu\text{M}$  sodium salts. All other chemicals were obtained from commercial sources.

### MelB overproduction and purification

As described previously (32), a  $\text{RecA}^-$  derivative of *E. coli* DW2 ( $\Delta\text{mel} \Delta\text{lacZY}$ ) was transformed with pK95 $\Delta\text{AHB}$  plasmid to overexpress a wild-type His-tagged MelB (8). Transformed cells were grown at  $30^\circ\text{C}$  in 200 L of M9 medium supplemented with appropriate carbon sources and ampicillin (100 mg/ml) at the Centre de Fermentation, Centre National de la Recherche Scientifique (Marseille, France), and used to prepare inverted membrane vesicles (IMVs), by means of a French press (American Instrument Exchange, Haverhill, MA). Purification of the His-tagged MelB was essentially carried out as previously described (33).

### Preparation of MelB proteoliposomes

Purified MelB protein (1–1.5 mg/ml) solubilized in dodecyl maltoside (DM, 0.1% w/v) was mixed with *E. coli* lipids to give a protein/lipid ratio of 1:2 (w/w). DM was removed by overnight adsorption in SM-2 Bio-Beads (Bio-Rad) at  $4^\circ\text{C}$  (34). The proteoliposomes were then subjected to repeated freeze/thaw-sonication-wash cycles in nominally  $\text{Na}^+$ -free buffer (20 mM MES, 10 mM KCl, pH 6.6) to eliminate the NaCl from the sample and medium.

### MelB activity and protein assays

MelB activity in proteoliposomes or in the solubilized state was assessed by measuring [ $\alpha$ - $^3$ H] NPG binding activity as previously described (2). The protein concentration was assayed using serum bovine albumin as standard.

### ATR-FTIR measurements

Twenty-microliter aliquots of MelB proteoliposome suspension (155  $\mu\text{g}$  protein), previously incubated with the desired substrates and briefly sonicated, were spread homogeneously on the upper surface of a germanium

crystal ( $50 \times 10 \times 2 \text{ mm}$ , yielding 12 internal reflections at the sample side, Harrick, Ossining, NY) under a stream of nitrogen. The film was left at room temperature for  $>1 \text{ h}$  and the ATR device was covered with a stainless steel plate containing a gas input and output. All the spectra were recorded at  $20^\circ\text{C}$  with a thermostatic circulatory bath connected to the stainless plate, using a FT6000 Bio-Rad spectrometer equipped with a cooled liquid nitrogen mercury-cadmium-telluride detector (Kolmar Technologies, Newburyport, MA) at a nominal resolution of  $2 \text{ cm}^{-1}$ .

Orientation experiments of MelB in  $\text{H}_2\text{O}$  (before H/D exchange) were performed with the incident light polarized parallel and perpendicular with respect to the incidence plane with a KRS polarizer (Perkin-Elmer, Wellesley, MA). Polarized spectra were also acquired after H/D exchange. H/D exchange experiments were performed basically as described in detail elsewhere (15). A film was exposed to a continuous flow of  $\text{N}_2$  gas (220 ml/min), saturated with  $^2\text{H}_2\text{O}$  by flushing it through nine closed 10-ml vials containing 2–3 ml of  $^2\text{H}_2\text{O}$  each. After 900 min exposure, the film was dried by flushing  $\text{N}_2$ , and polarized absorbance spectra were then acquired. Linear dichroism spectra were then computed by subtracting the perpendicular polarized spectrum from the parallel polarized spectrum with a suitable scaling factor (see below). Each polarized spectrum is the average of 2000 interferograms at a spectral resolution of  $2 \text{ cm}^{-1}$ .

Four film conditions were analyzed: either no substrates were added ( $\text{H}^+$  is the only substrate) or samples were equilibrated for 12 h in the presence of 10 mM NaCl, 5 mM melibiose, or 5 mM melibiose plus 10 mM NaCl. All samples were sonicated shortly before use. In the rehydrated state (where H/D exchange was performed) and the dry state (where polarized spectra were recorded), the relative concentrations of the media salt and sugar constituents increase to their respective solution states, ensuring full occupancy of the protein binding sites.

### Analysis of MelB secondary-structure tilting using ATR-FTIR dichroism spectra

Gaining information about structure orientation relies on the dependency of the infrared intensity on the angle between the transition dipole moment of a vibrational mode and the exciting electric field. By using polarized infrared radiation, these angles can be calculated, provided that the sample has a sufficient degree of orientation. The experimental quantity used to compute the structure orientation is the dichroic ratio  $R^{\text{ATR}}$ , defined as the ratio between the absorption of light polarized parallel and that polarized perpendicular to the incidence plane (35). On the other hand, the orientation is quantified by the experimentally accessible order parameter,  $S_{\text{experimental}}(\varphi)$ , which is a time- and space-averaged function of the angle between the transition dipole moment of a vibration and the  $z$  axis.

A perfect alignment with the  $z$  axis ( $\varphi = 0^\circ$  or  $\varphi = 180^\circ$ ) gives  $S_{\text{experimental}} = 1$ , whereas a perfect alignment in the  $xy$  plane ( $\varphi = 90^\circ$ ) yields  $S_{\text{experimental}} = -1/2$ . Alignment at the magic angle ( $\varphi = 54.7^\circ$ ) gives  $S_{\text{experimental}} = 0$ , but an isotropic arrangement (there is no order) also gives an  $S_{\text{experimental}}$  equal to zero. Hence, these conditions cannot be discerned based solely on the  $S_{\text{experimental}}$  value. In fact,  $S_{\text{experimental}}$  can be considered as a segmental parameter, since it is the product of three suborder parameters: the order parameter of the absorbing vibration dipole with respect to a molecular or a structure axis,  $S_{\text{dipole}}(\alpha)$ ; the order parameter of the molecular or structural axis with respect to the bilayer normal,  $S_{\text{structure}}(\beta)$ ; and the order parameter of the bilayer normal with respect to the  $z$  axis,  $S_{\text{membrane}}(\gamma)$  (see Fig. 23 in Goormaghtigh et al. (35)):

$$S_{\text{experimental}}(\varphi) = S_{\text{membrane}}(\gamma) \times S_{\text{structure}}(\beta) \times S_{\text{dipole}}(\alpha). \quad (1)$$

$S_{\text{experimental}}$  can be estimated from polarized ATR experiments, using the relation (35)

$$R^{\text{ATR}} = \frac{\int Abs_{\parallel}}{\int Abs_{\perp}} = \frac{\langle E_x^2 \rangle}{\langle E_y^2 \rangle} + \frac{\langle E_z^2 \rangle}{\langle E_y^2 \rangle} \times \frac{2S_{\text{experimental}} + 1}{1 - S_{\text{experimental}}}, \quad (2)$$

where  $R^{\text{ATR}}$  is the experimentally determined dichroic ratio from an ATR experiment and  $\langle E_x^2 \rangle$ ,  $\langle E_y^2 \rangle$ , and  $\langle E_z^2 \rangle$  are the time-averaged electromagnetic field intensities along the  $x$ ,  $y$ , and  $z$  directions of the ATR crystal (35). Their accurate values cannot usually be computed, since they depend on some parameters difficult to know, as the complex refraction index of the sample or the film thickness. However, if some experimental conditions are preserved, the thick film hypothesis provides a simple way for a simple calculation of  $\langle E_x^2 \rangle$ ,  $\langle E_y^2 \rangle$ , and  $\langle E_z^2 \rangle$  values (18,35,36). This hypothesis assumes that the sample is a weak absorber and the film thickness is much larger than the evanescent wave penetration for all the studied wavenumbers. These assumptions hold in our experiments, since all proteoliposome vibrations can be considered as weak absorbers and, even if completely dried, the amount of sample allows for at least 0.93- $\mu\text{m}$  film thickness, which is higher than the evanescent wave penetration above 700  $\text{cm}^{-1}$ . Therefore, in the following we will use the thick-film hypothesis.

In general, the polarizers are not able to polarize the infrared light to 100%, i.e., they have some leak. We determined our polarization leak to be 0.024 from the experimentally determined isotropic ratio of 1.93 (see below). This leaking is consistent with typical reported values (37). If we take into account the estimated leak of the polarizer in Eq. 2, then

$$R^{\text{ATR}} = \frac{\int Abs_{\parallel}}{\int Abs_{\perp}} \approx \frac{(1-f) \int Abs_{\parallel}^{\text{ideal}} + f \int Abs_{\perp}^{\text{ideal}}}{(1-f) \int Abs_{\perp}^{\text{ideal}} + f \int Abs_{\parallel}^{\text{ideal}}} = \frac{[\langle E_x^2 \rangle(1 - S_{\text{experimental}}) + \langle E_z^2 \rangle(2S_{\text{experimental}} + 1)](1-f) + \langle E_y^2 \rangle(1 - S_{\text{experimental}})f}{\langle E_y^2 \rangle(1 - S_{\text{experimental}})(1-f) + [\langle E_x^2 \rangle(1 - S_{\text{experimental}}) + \langle E_z^2 \rangle(2S_{\text{experimental}} + 1)]f}. \quad (3)$$

Using the expressions for  $\langle E_x^2 \rangle$ ,  $\langle E_y^2 \rangle$ , and  $\langle E_z^2 \rangle$  under the thick film/weak absorber approximation (35), taking into account our experimental setup (angle of the infrared (IR) beam with the ATR surface normal  $\theta = 45^\circ$ , refractive index of the ATR internal reflection element  $n_1 = 4$ ), and simplifying, we obtain

$$S_{\text{experimental}} = \frac{(n_2^2 - 16)(128R^{\text{ATR}} - 247)}{n_2^2(128R^{\text{ATR}} - 247) - 1904(R^{\text{ATR}} + 1)}, \quad (4)$$

where  $n_2$  is the real refraction index of the sample, which is usually assumed to be close to 1.5 for proteoliposomes (35). From Eq. 4, we obtain an experimental order parameter using the  $R^{\text{ATR}}$  of the polarized infrared spectra, which in turn can be used in Eq. 1 to estimate a structure order parameter or any other order parameter. To calculate  $R^{\text{ATR}}$  for different vibrations, the following integration intervals were used: 3400–3200  $\text{cm}^{-1}$  (amide A); 1696–1605  $\text{cm}^{-1}$  (amide I and amide I-I'); 1750–1720  $\text{cm}^{-1}$  (lipid C=O stretching); and 2931–2912  $\text{cm}^{-1}$  for lipid  $\text{CH}_2$  asymmetrical stretching. With the aim of obtaining a  $R^{\text{ATR}}$  value for the amide I as accurately as possible, the amino acid side-chain contribution was subtracted before  $R^{\text{ATR}}$  was calculated. The amino acid side-chain contribution synthetic data in  $\text{H}_2\text{O}$  and  $\text{D}_2\text{O}$  were constructed from published data (38,39). For the lipid  $\text{CH}_2$  stretching, the areas used to calculate  $R^{\text{ATR}}$  were measured from the Fourier deconvoluted spectra, to reduce the contribution of nearby bands. All errors reported in Table 1 and throughout this article correspond to standard errors obtained from at least three repetitions.

It should be taken into account that when  $R^{\text{ATR}}$  is obtained from a band area with overlapped subcomponent bands (such as the amide I band of proteins), the experimental order parameter should better be interpreted as a global experimental order parameter,  $S_{\text{experimental}}^{\text{global}}$ . Once the order-parameter properties are given, the global order parameter will be an average of the subcomponent bands order parameters, weighed by their area (35). If the amide I subcomponents are assembled in two groups, one will correspond to the averaged order parameter of the amide groups located in transmem-

brane helices, with a fractional contribution to the amide I total area,  $m$ ; the other will correspond to the averaged order parameter of the remainder of structures contributing to the amide I:  $S_{\text{experimental}}^{\text{global}} = m_{\text{helix}} S_{\text{experimental}}^{\text{helix}} + (1 - m_{\text{helix}}) S_{\text{experimental}}^{\text{rest}}$ . In our case, the amide I area fraction corresponding to transmembrane helices has been estimated to be  $\sim 0.5$  (8). If we consider the averaged order parameter of amide I vibrations of the structures other than helices to be close to zero, then  $S_{\text{experimental}}^{\text{helix}} = S_{\text{experimental}}^{\text{global}}/m$ . Once  $S_{\text{experimental}}^{\text{helix}}$  is obtained, the order parameter of the molecular axis of the helices,  $S_{\text{structure}}^{\text{helix}}$ , can be determined as

$$S_{\text{structure}}^{\text{helix}} = \frac{S_{\text{experimental}}^{\text{helix}}}{S_{\text{membrane}} \times S_{\text{dipole}}^{\text{helix}}} = \frac{S_{\text{experimental}}^{\text{global}}}{m \times S_{\text{membrane}} \times S_{\text{dipole}}^{\text{helix}}}. \quad (5)$$

This equation allows us to estimate the averaged order parameter (and tilt angle) of the transmembrane helices from the amide I band. The above assumption, considering the order parameter of nonhelical structures to be zero for amide I vibrations, does not mean that they have no orientation. Instead, we are assuming that the amide I vibration (roughly equivalent to the direction of the amide C=O bond) in nonhelical structures will have no net preferential orientation, and so their order parameter will average out to

nearly zero. This assumption is strongly supported when the deconvoluted amide I absorbance and the linear dichroism spectra are compared (compare Dave et al. (8), Fig. 5, with this article (see Fig. 3)). In the absorbance spectrum, many bands are observed, and only about half of the area is assignable to transmembrane helices (8). In contrast, fewer bands are observed in the linear dichroism spectrum, and  $>80$ – $85\%$  of its area is assignable to transmembrane helices. The  $S_{\text{experimental}}^{\text{helix}}$  value obtained from Eq. 5 will be subjected to some uncertainties, related to the use of the  $m$  value and the assumption that  $S_{\text{experimental}}^{\text{rest}} \approx 0$ . However, even if there is a 20% relative increase in  $m$  (from  $m = 0.5$  to  $m = 0.6$ ), the calculated tilt angle for the transmembrane helices will increase only  $\sim 3$ – $5^\circ$ . Moreover, since these corrections are the same for all the substrate conditions, the tilt differences between substrate conditions will remain basically unaffected, as well as our conclusions about substrate effects on transmembrane helix orientation.

Once the  $S_{\text{structure}}$  value is determined for the different molecular axes, we obtain the corresponding angle  $\beta$ :

$$\beta = \arccos \left[ \sqrt{\frac{2S_{\text{structure}} + 1}{3}} \right]. \quad (6)$$

This equation assumes a uniaxial orientation for the molecular axis. The molecular angle  $\beta$  will henceforth be referred to in this article as the average tilt angle of the structure.

## Reliability on structure tilting changes

Knowing  $S_{\text{experimental}}$ ,  $S_{\text{dipole}}$ , and  $S_{\text{membrane}}$  and applying Eq. 1 (or the modified Eq. 5 for the amide I) enables the computation of the molecular structures order parameter. The actual values of  $S_{\text{dipole}}$  for the amide vibrations in  $\alpha$ -helices have been controversial, with important discrepancies

**TABLE 1** Order-parameter values for vibrations and tilt angles of some molecular structures in the MelB proteoliposomes for different substrate conditions

	H <sup>+</sup>	Na <sup>+</sup>	H <sup>+</sup> -mel	Na <sup>+</sup> -mel
Order parameters				
Protein amide I*	0.15 ± 0.01	0.10 ± 0.01	0.15 ± 0.01	0.13 ± 0.01
H/D resistant protein amide A <sup>†</sup>	0.44 ± 0.03	0.32 ± 0.01	0.40 ± 0.01	0.40 ± 0.04
Lipid C=O	−0.26 ± 0.01	−0.24 ± 0.01	−0.25 ± 0.01	−0.24 ± 0.01
Lipid C-H <sup>‡</sup>	−0.33 ± 0.01	−0.32 ± 0.01	−0.33 ± 0.01	−0.32 ± 0.01
Lipid C-H <sup>§</sup>	−0.27	−0.27	−0.27	−0.27
Tilt angles <sup>¶</sup>				
Helices	26 ± 2	36 ± 1	25 ± 1	30 ± 1
H/D resistant helices	28 ± 2	36 ± 1	31 ± 1	31 ± 3

All order parameters and tilt angles are given with respect to the lipid membrane normal.

\*Calculated from the amide I in both H<sub>2</sub>O and D<sub>2</sub>O. Roughly corresponding to the order parameter of the amide C=O bond.

<sup>†</sup>Calculated from the amide A in D<sub>2</sub>O. Corresponds to the order parameter of the amide N-H bond.

<sup>‡</sup>From the lipid asymmetric CH<sub>2</sub> stretching.

<sup>§</sup>From the lipid symmetric CH<sub>2</sub> stretching, used as an fixed internal standard to correct for differences in membrane orientation over the ATR crystal (see Materials & Methods).

<sup>¶</sup>Details about the assumptions involved can be found in Materials & Methods and Results.

between authors (see, e.g., Table IIIB in Goormaghtigh et al. (40)). Recently, Marsh et al. experimentally determined the  $S_{\text{dipole}}$  of the amide vibrations (amides A, I, and II) using an  $\alpha$ -helix model polypeptide (41). The  $S_{\text{dipole}}$  values obtained were in very good agreement with values estimated later from the correlation of orientation experiments of bacteriorhodopsin with its x-ray crystallographic structure (42). Following Marsh et al., we used  $S_{\text{dipole}}$  0.65 for the amide A vibration mode, 0.43 for the amide I, and −0.37 for the amide II (41). It should be noted that the use of different values for  $S_{\text{dipole}}$  would affect the tilt angles estimated in this work, but the substrate-induced tilt angle differences should remain basically unaffected. A more serious problem could be the possibility that helices having different structures also could have a different  $S_{\text{dipole}}$ , as suggested by a theoretical work by Reisdorf and Krimm on  $\alpha_1$  and  $\alpha_{II}$  helices (43). In such a case, a difference in  $S_{\text{dipole}}$  between two types of helices could be artifactually interpreted as a difference in their  $S_{\text{structure}}$ , and so as having a different tilt angle. However, it should be pointed out that the relative differences in  $S_{\text{dipole}}$  calculated by Reisdorf and Krimm for  $\alpha_1$  and  $\alpha_{II}$  helices were only in the order of 10% despite their significant structural differences, and no significant difference in  $S_{\text{dipole}}$  as a function of the helix length was observed. Moreover, we should note that the work of Marsh and Pali suggests no significant differences in  $S_{\text{dipole}}$  for the different helices of bacteriorhodopsin (42). Therefore, we will make the reasonable assumption that all the helices in MelB have the same  $S_{\text{dipole}}$  independent of their actual conformation, length, etc. For phospholipids, the order parameter of the CH<sub>2</sub> stretching and scissoring with respect to the lipid chain axis, is  $S_{\text{dipole}} = -0.5$  (44). On the other hand, to determine accurate values and be able to reliably detect changes of  $S_{\text{structure}}$ , it is important to have a good estimate of  $S_{\text{membrane}}$  for each sample. For this purpose, it would be interesting to know  $S_{\text{structure}}$  for *E. coli* lipids to obtain  $S_{\text{membrane}}$ , since the  $S_{\text{structure}}$  for *E. coli* lipids is little affected by the presence of the protein and substrates. This last conclusion is based on the invariance of the lipid CH<sub>2</sub> stretching band maxima under different MelB substrate conditions, which is a very sensitive marker to detect changes on the lipid  $S_{\text{structure}}$  (45).

Unfortunately, no  $S_{\text{structure}}$  values for *E. coli* lipids are found in the literature, so we estimated it in the following way. The order parameter for a lipid chain of DMPC in the crystalline state is  $S_{\text{structure}} \approx 0.93$  (46). We obtained  $S_{\text{experimental}} = -0.356 \pm 0.03$  for the CH<sub>2</sub> stretching of DMPC in a dry film (pseudocrystalline state), which is roughly equivalent to a crystalline state. Applying Eq. 1, we obtained  $S_{\text{membrane}} = 0.77 \pm 0.04$  for the DMPC films. We assumed that the same  $S_{\text{membrane}}$  value will be valid for other lipid films prepared in the same way, such as *E. coli* lipids. In this way, we determined, for *E. coli* lipids,  $S_{\text{structure}} \approx 0.54$ . Using this value, together with  $S_{\text{experimental}}$  and  $S_{\text{dipole}}$  of the lipid CH<sub>2</sub> stretching we were able to obtain an accurate estimation of  $S_{\text{membrane}}$  for each MelB film.

## Linear dichroism spectra

Linear dichroism spectra were obtained as  $\text{Abs}_{\parallel} - R^{\text{iso}} \times \text{Abs}_{\perp}$ . The  $R^{\text{iso}}$  is the dichroic ratio for a band with an order parameter equal to zero (not oriented or tilted at the magic angle). We experimentally determined this ratio from the dichroic ratio of pure liquid water (a nonordered substance), as 1.93. The band intensity in a linear dichroism spectra is proportional to its absorbance  $\times$  its experimental order parameter. Therefore, the dichroism spectra properties are 1), bands showing  $S_{\text{experimental}} \approx 0$  cancelled; 2), bands with  $S_{\text{experimental}} > 0$  appearing as positive bands; and 3), bands with  $S_{\text{experimental}} < 0$  appearing as negative bands. Due to the disappearance of bands with  $S_{\text{experimental}} \approx 0$ , the resolution of bands arising from oriented structures is enhanced. Furthermore, the resolution of these oriented bands is further enhanced by using Fourier self-deconvolution (FSD) (47,48), using slightly different parameters depending on the region. In the 1800–1580 cm<sup>−1</sup> region, FSD was applied with a full width at half-height ( $\gamma$ ) of 18 cm<sup>−1</sup>, a band-narrowing factor ( $k$ ) of 2.5 and a Bessel filter. In the 1590–1450 cm<sup>−1</sup> region, where some narrow bands are present, FSD was performed with  $\gamma = 18$  cm<sup>−1</sup>,  $k = 2.2$ , and a Gaussian filter. For the purpose of quantifying amide I and amide I-I', FSD spectra were curve-fitted using Fourier deconvoluted Lorentzian bands with probabilistic constraints in the positions ( $\pm 2$  cm<sup>−1</sup>) and bandwidths ( $\pm 8$  cm<sup>−1</sup>), as described in detail elsewhere (49,50).

## Obtaining the order-parameter spectra

The structural order parameter for the amide I as a function of wavenumber was obtained in the following way. First, deconvoluted parallel and perpendicular polarized absorbance spectra (with the amino acid side chain contribution subtracted) were rationed, giving  $R^{\text{ATR}}$  versus wavenumber. This dichroic ratio spectrum was converted to a structural order parameter spectrum by applying Eq. 4 and then Eq. 1, and using the  $S_{\text{dipole}}$  corresponding to the amide I of an  $\alpha$ -helix. FSD ( $\gamma = 18$  cm<sup>−1</sup>,  $k = 2.2$ , and a Gaussian filter) was used to decrease the band overlapping, as well as improving the accuracy of the obtained order parameters. Further details about the interpretation of order parameter spectra are given as Supplementary Material.

## RESULTS

To obtain information on MelB secondary-structure orientation and analyze the changes in tilting caused by substrate

interaction, experiments were carried out on dry proteoliposomes prepared with purified MelB. Proteoliposome samples were fully equilibrated with media differing in substrate composition, favoring the formation of different MelB intermediates of the catalytic cycle, namely  $H^+ \cdot \text{MelB}$ ,  $Na^+ \cdot \text{MelB}$ ,  $H^+ \cdot \text{mel} \cdot \text{MelB}$ , and  $Na^+ \cdot \text{mel} \cdot \text{MelB}$  complexes.

As detailed in Materials and Methods, the order parameters directly obtained from the experimental data are relative to the ATR crystal normal, whereas the order parameter and especially the tilting of any of the MelB secondary-structure components are of interest if expressed with respect to the bilayer normal. To convert the reference frame, the experimental order parameter,  $S_{\text{experimental}}$ , must be corrected by the order parameter of the mosaic spread,  $S_{\text{membrane}}$  (order parameter of the membrane normal with respect to the ATR crystal normal). For each sample,  $S_{\text{membrane}}$  was determined from Eq. 1, using the  $S_{\text{experimental}}$  of the  $CH_2$  symmetric stretching from the lipids, and  $S_{\text{structure}} = 0.54$  and  $S_{\text{dipole}} = -0.5$  (see Materials and Methods for details). The average mosaic spread order parameter obtained for the MelB proteoliposomes under all conditions examined was  $0.69 \pm 0.09$ . By comparing this value with that obtained for pure DMPC lipid ( $S_{\text{membrane}} \approx 0.77$ ) it can be inferred that the presence of the protein reduces, to some extent, the film order over the ATR crystal, as we expected. Using the  $S_{\text{membrane}}$  value estimated for each MelB film, we could express the vibration order parameters and tilt angles of molecular structures in MelB with respect to the lipid bilayer normal, no longer to the ATR crystal normal.

### Polarized and dichroism spectra of MelB in $H_2O$ medium without substrates

Preliminary information on the presence of oriented secondary structures in MelB proteoliposomes equilibrated with  $Na^+$ - and sugar-free aqueous buffer ( $H^+ \cdot \text{MelB}$  species) was obtained by recording absorption spectra with perpendicular and parallel polarized IR radiations (Fig. 1 *a*). The amide I, II, and A maxima wavenumbers (1658, 1544, and 3298  $cm^{-1}$ , respectively) are typical of membrane proteins with a main contribution of helical structures (40). However, the small shoulder in the amide A at 3200  $cm^{-1}$ , as well as a previous quantitative analysis of the amide I (8), indicate some significant contribution of  $\beta$ -sheet structures, as well as turns.

The dichroism spectrum, enlarged in Fig. 1 *b*, was calculated by subtracting the perpendicular ( $\perp$ ) spectrum from the parallel ( $\parallel$ ) spectrum, corrected by a subtraction factor of 1.93 (see Materials and Methods). From the negative sign of the  $CH_2$  and  $C=O$  lipid bands in the dichroism spectrum, we can deduce that the film is correctly oriented over the ATR crystal. The presence of signals at wavenumber intervals corresponding to the amide A (3296  $cm^{-1}$ , positive), amide I (1660.5  $cm^{-1}$ , positive), and amide II (1543.2  $cm^{-1}$ , negative) bands clearly indicates that MelB contains a significant amount of oriented secondary-structural components. More-

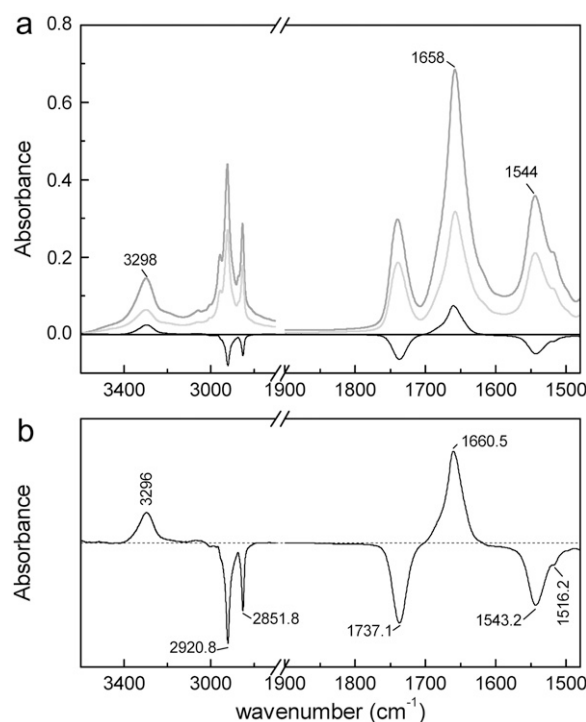


FIGURE 1 (a) Absorbance spectra of a dry film of MelB proteoliposomes in buffer (20 mM MES and 10 mM KCl, pH 6.6; no substrates added) for perpendicular (gray) and parallel (light gray) polarized light, together with the corresponding linear dichroism spectrum (black line). (b) Expanded linear dichroism spectrum.

over, the amide A, I, and II wavenumber maxima are all assignable to  $\alpha$ -helices, indicating that the major part of the oriented structures corresponds essentially to transmembrane helices. One may note the coherence of the amide A, I, and II signal signs, because the major vibration contributing to the amide I normal mode is due to the  $C=O$  stretching, which is parallel to  $NH$  stretching (amide A vibration), and perpendicular to  $NH$  bending, the major contribution to the amide II normal mode.

### Polarized and dichroism spectra of MelB in $D_2O$ medium without substrates

Similar analysis was carried out on MelB film previously fluxed for 900 min with  $D_2O$  vapor (Fig. 2 *a*). In this situation, and in agreement with previous measurements (15), 40% of the amide group hydrogens appeared to be resistant to the  $H/D$  exchange, as determined from the amide II area decrease and after amino acid side-chain corrections. From the amide A and II bands maxima (3302 and 1543  $cm^{-1}$ , respectively) and the symmetric shape of the amide A, we can deduce that the  $H/D$ -resistant amide hydrogens are mostly located in helices.

The enlarged dichroism spectrum is presented in Fig. 2 *b*. Amide A and amide II correspond to  $H/D$ -resistant structures,

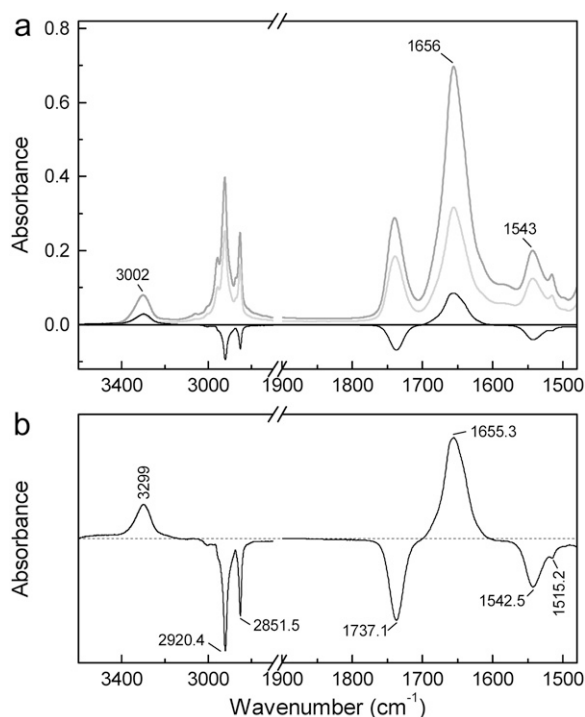


FIGURE 2 (a) Absorbance spectra of a dry film of MelB proteoliposomes in buffer (20 mM MES and 10 mM KCl, pH 6.6; no substrates added) for perpendicular (gray) and parallel (light gray) polarized light after 900 min incubation in D<sub>2</sub>O, together with the corresponding linear dichroism spectrum (black line). (b) Expanded linear dichroism spectrum.

which in addition are oriented. From the signal intensity and signs arising in the amide A and II bands (maxima at 3299.0 and 1542.5 cm<sup>-1</sup>, respectively) and the symmetric shape of the amide A, we can deduce that the oriented structures resistant to H/D exchange are most likely transmembrane helices (52), although there exists a certain amount of other oriented structures. On the other hand, the analysis of amide A' and amide II' (potential characterizers of the structure adopted by the H/D-exchanged amide hydrogens) presented unworkable difficulties. The amide A' gave two bands in the dichroism spectrum, at 2590 and 2470 cm<sup>-1</sup> (data not shown). However, since no amide A' assignments are currently available in the literature it was not possible to assign the corresponding secondary structures. The amide II', totally overlapped with a lipid band at 1460 cm<sup>-1</sup>, is therefore not experimentally observable. Finally, the analysis and interpretation of the highly overlapped amide I and I' signals (amide I-I') is challenging (35,53), since it represents a mixture of the nonexchanged and exchanged oriented structures. This problem was solved by combining deconvolution and curve-fitting, as described below.

### Deconvoluted dichroism spectra in H<sub>2</sub>O or D<sub>2</sub>O medium without substrates

To improve the band resolution and to identify the oriented secondary-structural components, the MelB dichroism spec-

tra were deconvoluted. In Fig. 3, the black trace shows a deconvoluted spectrum in H<sub>2</sub>O medium devoid of substrates (i.e., in the presence of H<sup>+</sup>). The most salient feature is the amide I showing an intense maximum at 1661.8 cm<sup>-1</sup> and a shoulder at 1654.0 cm<sup>-1</sup>. According to theoretical studies, an ideal transmembrane  $\alpha$ -helix gives two amide I vibrations with opposite orientation: a first main component parallel to the helix axis (A mode), and a second one, around five times less intense, perpendicular to the helix axis (E mode) (43). Hence, the band observed at 1654.0 cm<sup>-1</sup> cannot be assigned to the lower component of the band at 1661.8 cm<sup>-1</sup>, since both bands show positive intensity in the dichroism spectrum and comparable intensity in the absorbance spectrum (8). From this, we infer that the two individualized bands correspond to dichroic signals arising from two different types of transmembrane  $\alpha$ -helices. It is worth noting that their respective positions coincide with those of the previously assigned type I and type II  $\alpha$ -helices (8). For amide II, only one band at 1544.7 cm<sup>-1</sup> is observed, meaning a major contribution of oriented  $\alpha$ -helices (52). Further examination of Fig. 3 suggests that other MelB secondary-structure components may also display a significant degree of orientation. Thus, the bands at 1682.6 and 1673.7 cm<sup>-1</sup> may represent open loops, contributed by the first amino acids belonging to the loops linking the transmembrane segments, which can adopt reverse turns structure (54–56) and should show orientation. Also, the band at 1643.7 cm<sup>-1</sup> may correspond to  $3_{10}$  helices or to  $\alpha$ -helices with strong hydrogen bonds (8). This band could be assigned to oriented  $\beta$ -sheet structures if the corresponding amide II signal at  $\sim$ 1535 cm<sup>-1</sup> was observed (52). The assignment could still be possible if the amide II component was hidden by the more intense helix band, but the amide II envelope does not give a clear support for the presence of an unresolved component at  $\sim$ 1535 cm<sup>-1</sup>. Furthermore, no band is observed in the 1640–1620 cm<sup>-1</sup> amide I region, a typical interval for  $\beta$ -sheet absorption (52). Finally, typical bands from tyrosine and phenylalanine side chains appear at 1516.0 cm<sup>-1</sup> and 1497.0 cm<sup>-1</sup>, respectively (38). Not surprisingly, about half of the tyrosines (13 out of 25) and phenylalanines (19 out of 33) are presumably located in the transmembrane segments (5).

In Fig. 3, the gray trace shows the deconvoluted dichroism spectrum of MelB after equilibration with a D<sub>2</sub>O saturated atmosphere for 900 min. In addition to the three main bands observed in the amide I' region (1661.9, 1652.6, and 1641.3 cm<sup>-1</sup>), two smaller bands are observed in the spectrum at 1682.6 and 1673.1 cm<sup>-1</sup>. In the amide II, a single band at 1543.8 cm<sup>-1</sup> appears with a very similar shape to that in H<sub>2</sub>O although with lower intensity with respect to the amide I-I' area, which works as an internal standard. The band at 1661.9 cm<sup>-1</sup> in D<sub>2</sub>O should correspond to the nonexchanged structure at 1661.8 cm<sup>-1</sup> of amide I in H<sub>2</sub>O. As mentioned above, precise assignment of the other dichroic signals to a given secondary structure is often complicated by possible overlapping of exchanged (amide I') with nonexchanged

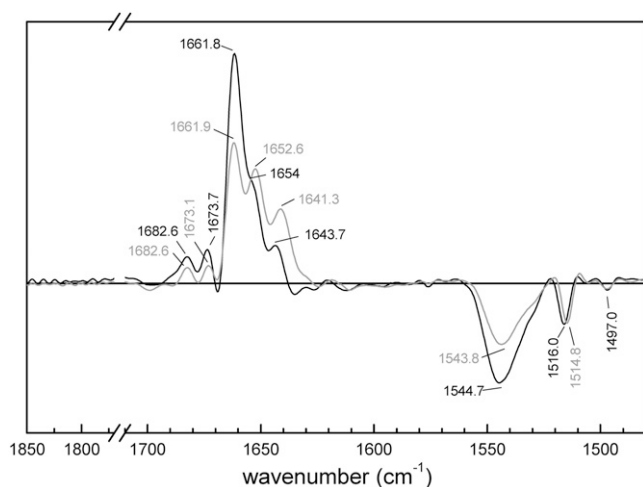


FIGURE 3 Deconvoluted linear dichroism spectra for a dry film of MelB proteoliposomes in buffer (20 mM MES and 10 mM KCl, pH 6.6; no substrates added) before (black) and after (gray) 900 min incubation in D<sub>2</sub>O.

(amide I) oriented structures. Thus, the 1652.6 cm<sup>-1</sup> band is probably due to contributions of the exchanged 1661.8 cm<sup>-1</sup> band and the nonexchanged 1654 cm<sup>-1</sup> band. Finally, the band at 1641.3 cm<sup>-1</sup> should be considered the result of an exchanged structure from the 1654 cm<sup>-1</sup> band plus the nonexchanged structure of the band at 1643.7 cm<sup>-1</sup>. In any instance, all these bands arise from  $\alpha$ -helical components. Taken together, the data in D<sub>2</sub>O medium therefore suggest that the predominant oriented structures in MelB are, in substance, mostly  $\alpha$ -helical. Incidentally, the band corresponding to tyrosine shifts from 1516.0 in H<sub>2</sub>O to 1514.8 cm<sup>-1</sup> in D<sub>2</sub>O, meaning that, despite their putative transmembrane location, the oriented tyrosines have the -OH group H/D-exchangeable (39).

### Effect of substrate on MelB secondary-structure tilting

Table 1 compares the order-parameter values of several protein and lipid vibrations with respect to the membrane normal, obtained as  $S_{\text{experimental}}/S_{\text{membrane}}$ , for MelB proteoliposomes equilibrated in the four different substrate conditions (H<sup>+</sup>, Na<sup>+</sup>, H<sup>+</sup>/melibiose, or Na<sup>+</sup>/melibiose). A comparison of the linear dichroism spectra under different substrate conditions is included as Supplementary Material (Fig. S1), in both H<sub>2</sub>O and D<sub>2</sub>O. This allows the interested reader to evaluate the substrate effects independently of the way the dichroic ratios are converted in structural order parameters and tilt angles.

One first important observation from Table 1 is that the lipid vibrations do not change their order significantly as a result of MelB substrate presence, in contrast to the changes clearly exhibited by the protein vibrations. The protein vibrations considered are the amide I (roughly equivalent to the direction of the amide C=O bonds), and the amide A after

900 min incubation in D<sub>2</sub>O (direction of the amide NH bond with hydrogens resistant to H/D exchange). These vibration order parameters are positive, which means that both the amide C=O and NH bonds show a preferential orientation toward the membrane normal in MelB, as expected for a protein containing a high percentage of amide groups involved in transmembrane helices.

To convert the order parameter of the amide I vibration into a structural order parameter of the helix axis, we need to take into account that we have other structures besides helices contributing to the amide I area, and thus to its experimental order parameter. Taking into account that the helix contribution to the amide I area is roughly 50%, we obtained the average tilt of the helices for each substrate condition, shown in Table 1. Note that in this calculus we also assumed that the order parameter of the amide I vibration averages to nearly zero, except for the amide groups located in helices. This assumption becomes justified after observing that although helix contribution to the amide I area in the absorbance spectra is ~50% (8), the area contribution in a linear dichroism is 85–90% (see Figs. 3 and 5). This implies that the average order parameter of an amide I vibration located in helices is more than eight times higher than the average order parameter of the amide I vibration located in other structures, i.e., the assumptions used to obtain the tilted angles in Table 1 are completely appropriate.

As deduced by the amide A and II maxima after D<sub>2</sub>O incubation (Fig. 2 *a*), the H/D-exchange-resistant amide NH bonds correspond to helices. Making the approximation that all the H/D-exchange-resistant amide groups contributing to the amide A area (3400–3200 cm<sup>-1</sup>) are located in helices, we estimated the tilt axis of the corresponding H/D-exchange-resistant helices. The calculated tilt angles for different substrates conditions are shown in Table 1.

Note that the estimated tilt angles for the MelB helices (from amide I in H<sub>2</sub>O and D<sub>2</sub>O) and those for the H/D-exchange-resistant helices (from amide A in D<sub>2</sub>O) are in good agreement, although the assumptions involved in the two cases were completely different. This confirms that H/D-exchange-resistant structures correspond mainly to transmembrane segments. More important, these values indicate that the transmembrane helices reorient upon substrate binding. The averaged angle calculated for H<sup>+</sup>·MelB and H<sup>+</sup>·mel-MelB complexes is ~26°, whereas the Na<sup>+</sup>·MelB complex has an angle around 36°, and finally, the Na<sup>+</sup>·mel-MelB complex shows an angle of 30°.

One of our aims was to get more specific information about the tilting of type I and II helices and their variation as a function of Na<sup>+</sup> and melibiose interaction. To this end, we evaluated the order parameter of the helices,  $S_{\text{structure}}$ , as a function of wavenumber (order-parameter spectrum) for amide I in the 1670–1650 cm<sup>-1</sup> region, since it includes the signals arising from the two helix types (Fig. 4). The deconvoluted polarized absorbance spectral data were used for the calculation to improve structural resolution. Details about

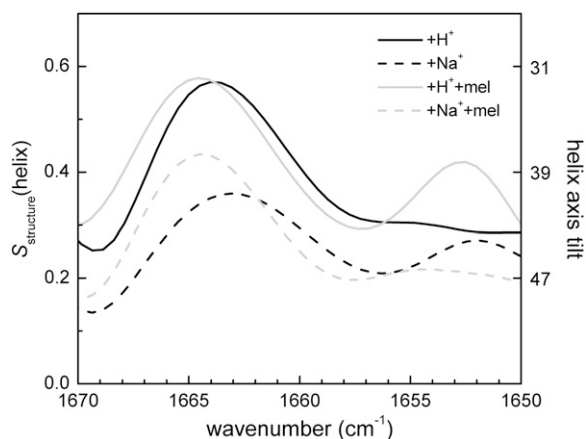


FIGURE 4 Estimated structure order parameter  $S_{\text{structure}}$  of helices in MelB versus the wavenumber for different substrate conditions: (black solid line) no added substrates, (black dashed line) addition of  $\text{Na}^+$ , (light gray solid line) addition of melibiose, and (light gray dashed line) addition of melibiose plus  $\text{Na}^+$ . The right axis shows the helix tilt angle ( $^\circ$ ) corresponding to the helix structure order parameter. The maxima values around 1664 and 1652  $\text{cm}^{-1}$  give an approximate lower bound to the tilt for helices I and II, respectively. For more detail, see text and Supplementary Material.

how Fig. 4 was obtained and how it should be interpreted are given as Supplementary Material (Fig. S2).

A clear maximum in the region at 1663–1665  $\text{cm}^{-1}$  appears in all substrate conditions, whereas a second, smaller maximum is only observed for two conditions at  $\sim 1653 \text{ cm}^{-1}$  (Fig. 4). It should first be noted that according to simulations using synthetic data, when two oriented (e.g., type I and II helices) and one nonoriented bands overlap, two maxima will usually appear in the order-parameter spectrum, at frequencies close (but not equal) to the absorbance maxima of each of the overlapping species (see Supplementary Material). Moreover, the order parameter at the band maxima will be representative of the oriented-band order parameter, giving in this case an indication of the tilt angle of the two types of helices. We should note, however, that these order parameters and tilt angles should be regarded as approximated lower bounds for the actual order parameters, and approximated higher bounds for the tilt angles.

Therefore, from the order-parameter spectra in Fig. 4, we can conclude that the tilting of both helical bands change

upon substrate binding (see the scale on the right). The highest maxima, at around 1664  $\text{cm}^{-1}$ , represent the less tilted helix with respect to the membrane normal, whereas the small maxima around 1653  $\text{cm}^{-1}$  correspond to more tilted helices. It is evident from Fig. 4 that the tilting of both types of helices is influenced by substrate interaction with the transporter.

### H/D exchange properties of $\alpha$ -helices type I and type II

Fig. 5 shows a representative curve-fitting of deconvoluted linear dichroism spectra of MelB without added substrates (Fig. 5 *a* in  $\text{H}_2\text{O}$ , and Fig. 5 *b* in  $\text{D}_2\text{O}$ ). Comparison of the fitted-band areas in  $\text{H}_2\text{O}$  and  $\text{D}_2\text{O}$  provides a means to evaluate the percentage of exchange of the structures giving rise to the two bands at 1661.8 and 1654.0  $\text{cm}^{-1}$  in  $\text{H}_2\text{O}$ , assigned to oriented transmembrane  $\alpha$ -helix types I and II, respectively (Table 2). H/D exchange for the 1661.7  $\text{cm}^{-1}$  band amounts to 36%, whereas that of the 1653.9  $\text{cm}^{-1}$  band is 53%. This calculation was repeated for diverse substrate conditions, and the results are collected in Table 2. The data suggest that either  $\text{Na}^+$  or melibiose binding to MelB decreases the H/D-exchange level of type I  $\alpha$ -helices, whereas the H/D exchange of type II  $\alpha$ -helices clearly decreases in the presence of melibiose.

## DISCUSSION

The results reported in this study provide evidence for the presence of an oriented structure that essentially corresponds to transmembrane helical domains, including, in particular, two previously identified populations of helices with different IR signatures ( $\sim 1661$  and  $\sim 1653 \text{ cm}^{-1}$ ). The data also suggest that both the average tilt angle and the orientation of each of these two particular tilted helices vary to a different extent according to the bound substrate(s).

### Oriented components in the MelB secondary structure

Analysis of MelB structure properties using IR polarized light indicates MelB contains a significant amount of tilted structural

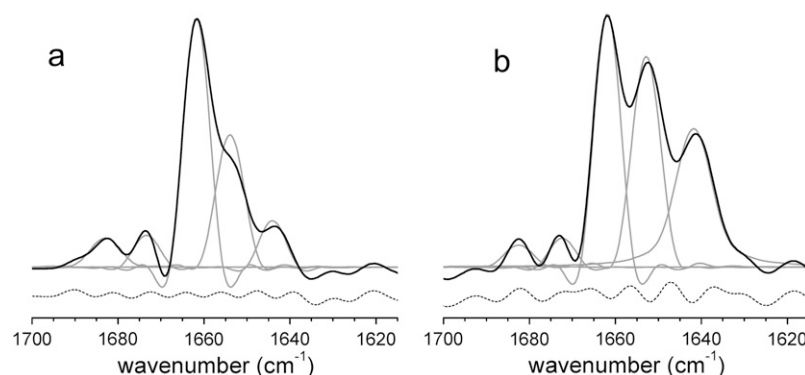


FIGURE 5 (a) Deconvoluted linear dichroism spectra for a dry film of MelB proteoliposomes in buffer (20 mM MES and 10 mM KCl, pH 6.6, with no added substrates (black solid line)). Curve-fitted bands are represented by gray lines. The fitting residual (short-dashed line) is shifted down for a clear visualization. (b) Same as *a*, but for a sample incubated for 900 min in  $\text{D}_2\text{O}$ .



**TABLE 2** Estimated positions and percentage of H/D exchange for helix types I and II in MelB from the curve-fitting of deconvoluted linear dichroism spectra

	H <sup>+</sup>	Na <sup>+</sup>	H <sup>+</sup> ·mel	Na <sup>+</sup> ·mel
Type I $\alpha$ -helices				
Position (cm <sup>-1</sup> )	1661.7	1661.5	1662.5	1663.1
H/D exchange (%)	36	22	19	22
Type II $\alpha$ -helices				
Position (cm <sup>-1</sup> )	1653.9	1653.4	1654.1	1655.0
H/D exchange (%)	53	46	37	37

The uncertainties in the H/D exchange quantification are approximately  $\pm 3$  for type I helices and  $\pm 5$  for type II helices.

components, forming an average angle of  $\sim 26^\circ$  with the membrane normal in a substrate-free medium. The dichroic bands in H<sub>2</sub>O or D<sub>2</sub>O media arise at  $\sim 1660$  and  $\sim 1655$  cm<sup>-1</sup> (amide I interval),  $\sim 1554$  cm<sup>-1</sup> (amide II interval), and  $\sim 3296$  and  $\sim 3299$  cm<sup>-1</sup> (amide A interval). These frequencies strongly suggest that the oriented structural components of MelB are essentially  $\alpha$ -helices. Furthermore, the deconvoluted linear dichroism spectra in H<sub>2</sub>O and D<sub>2</sub>O make it apparent that three individualized  $\alpha$ -helical populations, differing in their absorbance frequency ( $\sim 1661$ ,  $\sim 1653$ , and  $\sim 1644$  cm<sup>-1</sup>), contribute for most of the oriented protein amide I vibrations. Their positive sign in the linear dichroism spectra probes a preferential orientation in the direction of the membrane normal, strongly suggesting that these helical structures correspond to transmembrane domains. It is noteworthy that the existence of a significant number of tilted transmembrane domains in the MelB structure is strongly suggested by the projection map or the surface-rendered 3D map derived from 2D crystals (21). From the available resolved 3D structures, we learn that large proportions of helices are transmembrane and tilted. This general feature is found in 12 TM transporters belonging either to the MSF transporter family, namely, the LacY symporter (24), the oxalate OxIT, and the Glycerol-3-P GltP antiporters (23,26), to the Na<sup>+</sup>-H<sup>+</sup> antiporter family (NhaA, (27)), or, finally, to the neurotransmitter sodium symporter family (the Na<sup>+</sup>-leucine symporter LeuTaa, (28)).

The average angle value for the H<sup>+</sup>·MelB binary complex, estimated as  $26^\circ$ , is of the same order of magnitude as those reported for the H<sup>+</sup>·lacY symporter, the K<sup>+</sup> tetrameric channel of *Streptomyces lividans* ( $33^\circ$ ) (19,36), or the EmrE transporter ( $27^\circ$ ) (18). However, this comparison is only qualitative, as different factors can affect the average tilt angle value. For example, the lipid/protein ratio used for protein reconstitution, which is known to modify the average tilt angle (36), may be different in these reconstituted proteins. We should note that in all these works the orientation was estimated from dry films of the reconstituted proteins, and the possibility exists that the orientation (and the tilt changes on substrate binding) could be different in solution. However, it is important to keep in mind that water is present in dry films of proteins, and the hydration level may be not very different

from the hydration level in the protein crystals used in x-ray diffraction (35).

### Presence of two major populations of oriented transmembrane helices

Distinction between  $\alpha$ -helical signals at  $1661$  cm<sup>-1</sup> and at  $1653$  cm<sup>-1</sup> has already been inferred from the Fourier self-deconvoluted amide I (8) and ATR-FTIR difference spectra of MelB (9,16). The corresponding  $\alpha$ -helices were termed, respectively, types I and II. The detection of structural components at similar frequencies in the linear dichroism spectra clearly indicates that type I and II helices are oriented. They account for a large proportion of the MelB oriented structures. It is noteworthy that they differ by their orientation with respect to the membrane normal, type I helices being proportionally less tilted than type II helices. In addition, and as will be discussed in more detail below, the tilt of type I and type II helices are modified in a different manner by the substrates. The physicochemical factors leading to the different maximal absorption frequency at which the two helix types are recorded is a complex issue. Theoretical and experimental studies show that the amide I frequency for an  $\alpha$ -helix increases as the helix length decreases (57,58), decreases as the distortions and irregularities of the peptide bond dihedral angles increase (59), decreases as the helix hydrations increase (57), and increases as the peptide backbone  $i \rightarrow i + 4$  hydrogen bond length increases (60). For type I and II helices, one could invoke, for example, a difference in their length. If this were the case for MelB, the type I  $\alpha$ -helix signal ( $1661$  cm<sup>-1</sup>), or part of it, would arise from short helices, whereas the type II  $\alpha$ -helix signal ( $1653$  cm<sup>-1</sup>), or—again—part of it, would be generated by longer helices. However, one might expect that all these distortions of the canonical structure lead not only to a wavenumber maximum shift, but also to a change in the dipole orientation of the amide vibrations with respect to the helix axis. Therefore, it is conceptually possible that the differences in the dichroic ratio between type I and type II helices may originate not from tilt differences, but just from structural differences (43). However, we consider that reasonable differences in the dipole orientation caused by structural changes are insufficient to explain the observed dichroic differences (see Materials and Methods), and the observed differences in tilt between type I and type II helices are highly likely to be real. Nevertheless, it is important to recognize that structural changes leading to a change in the dipole orientation, but not affecting the helix tilt angle, can be misleadingly interpreted as a change in the helix tilt.

### Changes of helical tilting upon substrate interaction

Undoubtedly, a most interesting result is the observed change of tilting of the transmembrane  $\alpha$ -helical components resulting from the interaction of each cosubstrate with MelB.

These changes are not only substrate-specific, but are also different for each of the two major type I and II  $\alpha$ -helix components (Fig. 4). A general trend is that type I helices are systematically less tilted than type II helices in all MelB intermediates.

The average MelB tilt increases when shifting from the  $H^+$ -MelB to the  $Na^+$ -MelB intermediate. This is associated with a significant increase in the tilt of type I  $\alpha$ -helices, and not—or at most only modestly—in that of type II  $\alpha$ -helices (Fig. 4). Associating the quite selective  $Na^+$ -dependent change in tilt angle of type I  $\alpha$ -helices and the possibility that they might correspond to short MelB helices raises the interesting possibility of a structural analogy between the helices lining the  $Na^+$ -binding site of the  $Na^+$ - $H^+$  antiporter NhaA or of the  $Na^+$ -leucine LeuTaa symporter and that of MelB (27,28). It should also be stressed that MelB and NhaA have comparable projection maps at medium resolution, showing two regions with different shape and helical arrangements (21,22). It is striking that the determined  $Na^+$  binding site of NhaA and that predicted for MelB are both located in the corresponding compact regions. As learned from the NhaA and LeuTaa structures, the structural determinant of their respective  $Na^+$  binding site (Na1 site in the case of LeuTaa) includes, in addition to a long helix (HV for NhaA and H7 for LeuTaa), two helices (H4 and HXI for NhaA and HI and H6 for LeuTaa) interrupted by an unwound stretch, generating two sets of two short helices. The stretch flanking regions harbor the side-chain or main-chain carbonyls expected to be involved in  $Na^+$  coordination. From an IR spectroscopic point of view, the resulting inner or outer “half-helices” of any interrupted helix would behave as short helices. In view of the polarization data presented in this study, we hypothesize that one or more helices surrounding the ion binding site of MelB may also be interrupted helices. Tilting of some of these “half-helices” to accommodate the bound  $Na^+$  may lead to a reduction of their accessibility to the solvent, in agreement with our H/D measurements. Finally, based on biochemical evidence suggesting that helices I, II, and IV of the N-terminal half of MelB harbor Asp residues that are critical for  $Na^+$  recognition or binding (32,61), it is suggested that one (or more) of these particular helices may contain an unwound stretch. It should finally be recalled that  $Na^+$  cooperatively enhances the affinity of MelB for the sugar. One cannot a priori exclude that part of the tilt variations of type I  $\alpha$ -helices, and perhaps also of the modest tilt change of type II  $\alpha$ -helices, also reflect structural change involved in the cooperative effect. This would be consistent with a recent suggestion by Gouaux and co-workers (62) that ion binding to  $Na^+$ -coupled secondary transporters causes not only local but also long-distance conformational changes through the bending and twisting of the associated helical elements.

As illustrated in Table 1 and Fig. 4, interaction of melibiose with  $H^+$ -MelB does not modify the MelB average angle. This result is similar to that reported for the interaction of lactose with the  $H^+$ -coupled LacY transporter (36). However,

melibiose interaction with  $Na^+$ -MelB decreases the average angle. When MelB operates according to the  $H^+$ -coupled mode, melibiose significantly reduces the tilt of type II  $\alpha$ -helices, whereas no modification of tilting of type I  $\alpha$ -helices is observed. In the  $Na^+$ -coupled mode, melibiose binding accentuates tilting of type II  $\alpha$ -helices but reduces that of type I  $\alpha$ -helices. The large diversity in the observed effects could be due in large part to the utilization of a different coupling ion species for the sugar transport reaction. This diversity is most likely related to the large variations of kinetic constants (substrate affinity and rate of transport), cosubstrate binding, and cooperative properties existing between MelB functioning as  $H^+$ - or an  $Na^+$ -coupled sugar transporter (8,12,13,63). It is also significant that melibiose seems to affect the tilt of type II  $\alpha$ -helices proportionally more than that of type I  $\alpha$ -helices, whereas  $Na^+$  has the opposite impact on the tilt of these two helix populations, as discussed above. Finally, it must be stressed that the sugar-induced structural signals reported here may not be limited to the sugar-binding reaction, but may be mixed with (or replaced by) structural events associated with cooperative phenomena and/or to subsequent reorientation of the ternary MelB complex. With the sole information currently at hand, it appears unrealistic to look for any detailed structural explanation for the effect of the sugar on various helix tilt variations of MelB operating as  $H^+$ - and/or  $Na^+$ -coupled transport reaction. At a minimum, however, our study with polarized IR light provides additional support for a previous suggestion that conformational changes occur at different stages of the MelB transport cycle (12,13,16,63) and involve tilt variations of the two major helical populations of the transporter structure (9,17).

Clearly, improvement of the resolution of MelB structure will be essential to critically assess the validity of some of the proposed interpretations and/or speculations. It is also evident that examination of the fate of these signals in the many available mutants of MelB with altered substrate selectivity or with impaired coupling and/or translocation capacity should be useful to unravel the structure-function relationships of the  $Na^+$ -coupled MelB transporter in particular and of  $Na$ -coupled symporters in general.

## SUPPLEMENTARY MATERIAL

To view all of the supplemental files associated with this article, visit [www.biophysj.org](http://www.biophysj.org).

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